

# On some alternating series involving zeta and multiple zeta values

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*Journal of Mathematical Analysis and Applications*, **475** (2019)

**Abstract** In this article, we study a class of conditionally convergent alternating series including, as a special case, the famous series  $\sum_{n \geq 2} (-1)^n \frac{\zeta(n)}{n}$  which links Euler's constant  $\gamma$  to special values of the Riemann zeta function at positive integers. We give several new relations of the same kind. Among other things, we show the existence of a similar relation for the Apostol-Vu harmonic zeta function which have never been noticed before. We also highlight a deep connection with the Ramanujan summation of certain divergent series which originally motivated this work.

**Keywords** Riemann zeta function; harmonic zeta function; Stirling numbers of the first kind; Stirling numbers of the second kind; Bernoulli numbers; Bernoulli numbers of the second kind; harmonic numbers; multiple zeta values; Ramanujan summation of divergent series.

**Mathematics Subject Classification (2010)** 11B73; 11B75; 11M06; 11M32; 40G99; 41A58.

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# Introduction

The first part of this article is devoted to the study of the conditionally convergent alternating series  $\nu_k$  defined by

$$\nu_k := \sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{n+k},$$

where  $\zeta(s)$  is the Riemann zeta function and  $k$  denotes an integral parameter. By a classical result (cf. [9, p. 66], [16, p. 62]), it is well known that  $\nu_0$  is Euler's constant

$$\gamma = \lim_{n \rightarrow \infty} \left\{ \sum_{j=1}^n \frac{1}{j} - \ln n \right\} = 0.5772156649 \dots$$

This remarkable connection between  $\gamma$  and the special values at positive integers of the Riemann zeta function goes back to Euler's early works on harmonic series [14]. Less famous but yet fairly well-known (cf. [9, p. 93], [17, Eq. (5.1)], [18, Eq. (1.5)]) is the relation

$$\nu_1 = \frac{\gamma}{2} - \frac{1}{2} \ln 2\pi + 1$$

sometimes called Suryanarayana formula. Recently, Blagouchine [7, p. 413] gave a general expression of these series  $\nu_k$  in the case where  $k$  is a positive integer:

$$\begin{aligned} \nu_k &= \frac{\gamma}{2} - \frac{\ln 2\pi}{k+1} + \frac{1}{k} \\ &+ \sum_{r=1}^{\lfloor \frac{k}{2} \rfloor} (-1)^r \binom{k}{2r-1} \frac{(2r)!}{r(2\pi)^{2r}} \zeta'(2r) + \sum_{r=1}^{\lfloor \frac{k-1}{2} \rfloor} (-1)^r \binom{k}{2r} \frac{(2r)!}{2(2\pi)^{2r}} \zeta(2r+1). \end{aligned} \quad (1)$$

This formula seems quite cumbersome but can be much simplified using the functional equation of  $\zeta$ . After some elementary transformations, we show that equation (1) can be reduced to the following equivalent (but much more pleasant) expression:

$$\nu_k = \frac{\gamma}{k+1} - \frac{1}{2} \ln 2\pi + \sum_{j=1}^{k-1} (-1)^j \binom{k}{j} \zeta'(-j) + C_k, \quad (2)$$

where  $C_k$  is a rational number (see Proposition 1). Moreover, this expression allows to highlight a deep connection between  $\nu_{2k}$  and the sum (in the sense of the Ramanujan summation of divergent series) of the series  $\sum_{n \geq 1} n^{2k} H_n$ , where  $H_n$  is the  $n$ th harmonic number (see Remark 3).

Next, in a second part, we introduce a generalization of these series series  $\nu_k$  replacing the zeta values by certain multiple zeta values. A natural extension may

be defined as follows: for all integers  $k \geq -1$  and  $p \geq 0$ , we consider the class of series  $(\nu_{k,p})$  with

$$\nu_{k,p} := \sum_{n=2}^{\infty} \frac{(-1)^n}{n+k} \zeta(n, \underbrace{1, \dots, 1}_p),$$

where

$$\zeta(s_1, s_2, \dots, s_k) = \sum_{n_1 > n_2 > \dots > n_k \geq 1} \frac{1}{n_1^{s_1} n_2^{s_2} \dots n_k^{s_k}},$$

so that the previous series  $\nu_k$  become  $\nu_{k,0}$ . Then we establish (see Proposition 2) the following identity which is the main result of this work: we have

$$\nu_{k,p} = \sum_{n=1}^{\infty} \frac{|g_n^{(k+1)}|}{n^{p+1}}, \quad (3)$$

where the rational numbers  $g_n^{(k)}$  are defined by

$$g_n^{(k)} := \frac{1}{n!} \sum_{j=1}^n \frac{s(n,j)}{j+k} \quad (k \geq 0, n \geq 1), \quad (4)$$

where  $s(n,j)$  denotes the Stirling numbers of the first kind (comprehensive informations on the Stirling numbers of the first and the second kind may be found in [3, 7, 13, 16, 19]). In particular, for  $k=1$ ,  $g_n^{(1)}$  are nothing else than the *Bernoulli numbers of the second kind* [3, 13, 16]. One can prove easily (see Lemma 4) that  $g_n^{(k)} = (-1)^{n+1} |g_n^{(k)}|$ , so that the rational numbers  $g_n^{(k)}$  alternate in sign. As a special case of equation (3), we derive the following result:

$$\nu_{k-1} = \sum_{n=1}^{\infty} \frac{|g_n^{(k)}|}{n} \quad (k \geq 0). \quad (5)$$

In the case  $k=1$ , we recover the classical Mascheroni's series for  $\gamma$  (cf. [7, p. 406], [16, p. 280]):

$$\gamma = \frac{1}{2} + \frac{1}{24} + \frac{1}{72} + \frac{19}{2880} + \frac{3}{800} + \frac{863}{362880} + \dots$$

Another notable consequence of formula (3) is the deduction of this nice formula (see Corollary 1):

$$\sum_{n=2}^{\infty} (-1)^n \frac{\zeta_H(n)}{n} = \gamma_1 + \frac{1}{2} \gamma^2 + \frac{\pi^2}{12},$$

where  $\zeta_H(s)$  denotes the Apostol-Vu harmonic zeta function [4, 5, 6] and  $\gamma_1$  is the first Stieltjes constant [8, 9].

Finally, in the last section, we highlight a relation between the series  $\nu_k$ , the Stirling numbers of the second kind  $S(n,k)$ , and the shifted Mascheroni series  $\sigma_r$  whose study was the main subject of [13] (see Proposition 3 and Example 4).

# 1 The case of a positive integer

In this section, we focus on the case of a positive integer  $k$  and give two independent proofs of our formula (2). More precisely, we prove the following proposition:

**Proposition 1.** For any positive integer  $k$ , we have

$$\nu_k = \frac{\gamma}{k+1} - \frac{1}{2} \ln 2\pi + \sum_{j=1}^{k-1} (-1)^j \binom{k}{j} \zeta'(-j) + C_k$$

with

$$C_k = \frac{1}{k} + \sum_{j=0}^{k-1} \binom{k}{j} \frac{H_j B_{j+1}}{j+1}, \quad (6)$$

where  $H_n$  are the harmonic numbers,

$$H_0 = 0, \quad H_n := 1 + \frac{1}{2} + \cdots + \frac{1}{n} \quad (n \geq 1),$$

and  $B_n$  are the Bernoulli numbers defined by means of the exponential generating function

$$\frac{x}{e^x - 1} = \sum_{n=0}^{\infty} B_n \frac{x^n}{n!} \quad (|x| < 2\pi).$$

In particular,  $B_0 = 1$ ,  $B_1 = -1/2$ ,  $B_2 = 1/6$ ,  $B_{2r+1} = 0$  for  $r \geq 1$ .

*Proof.* We can quite easily derive (2) from (1). Differentiation of the functional equation [5, Eq. (25.4.2)]

$$\zeta(s) = 2(2\pi)^{s-1} \Gamma(1-s) \zeta(1-s) \sin \frac{\pi s}{2}$$

leads to the two relations

$$(-1)^r \frac{(2r)!}{2(2\pi)^{2r}} \zeta(2r+1) = \zeta'(-2r) \quad (r \geq 1),$$

and

$$(-1)^r \frac{(2r)!}{r(2\pi)^{2r}} \zeta'(2r) = -\zeta'(1-2r) + \frac{B_{2r}}{2r} (H_{2r-1} - \gamma - \ln 2\pi) \quad (r \geq 1).$$

Substituting these relations into (1) and grouping together the terms under the two symbols  $\Sigma$ , leads to the expression

$$\nu_k = \frac{\gamma}{k+1} - \frac{1}{2} \ln 2\pi + \sum_{j=1}^{k-1} (-1)^j \binom{k}{j} \zeta'(-j) + C_k,$$

where the rational constant  $C_k$  is given by equation (6). □

**Remark 1.** Formula (2) can also be rewritten as follows:

$$\nu_k = \frac{\gamma}{k+1} + \frac{1}{k} - \sum_{j=0}^{k-1} (-1)^j \binom{k}{j} L_j,$$

where

$$L_j := \frac{H_j B_{j+1}}{j+1} - \zeta'(-j) \quad (j \geq 0).$$

Furthermore, the numbers  $A_j := \exp(L_j)$  are known as the *Bendersky-Adamchik constants* [1, 2].

**Example 1.** For the first values of  $k$ , we have the following relations:

$$\begin{aligned} \nu_1 &= \frac{\gamma}{2} - \frac{1}{2} \ln 2\pi + 1, \\ \nu_2 &= \frac{\gamma}{3} - \frac{1}{2} \ln 2\pi - 2\zeta'(-1) + \frac{2}{3}, \\ \nu_3 &= \frac{\gamma}{4} - \frac{1}{2} \ln 2\pi - 3\zeta'(-1) + 3\zeta'(-2) + \frac{7}{12}, \\ \nu_4 &= \frac{\gamma}{5} - \frac{1}{2} \ln 2\pi - 4\zeta'(-1) + 6\zeta'(-2) - 4\zeta'(-3) + \frac{47}{90}, \\ \nu_5 &= \frac{\gamma}{6} - \frac{1}{2} \ln 2\pi - 5\zeta'(-1) + 10\zeta'(-2) - 10\zeta'(-3) + 5\zeta'(-4) + \frac{167}{360}. \end{aligned}$$

**Remark 2.** Starting from the Maclaurin series expansion [5, Eq. (25.8.5)]

$$\psi(x+1) + \gamma = \sum_{n=2}^{\infty} (-1)^n \zeta(n) x^{n-1} \quad (|x| < 1)$$

where  $\psi(x)$  denotes the digamma function (i.e. the logarithmic derivative of the  $\Gamma$ -function), and multiplying each side by  $x^k$  (with  $k \geq 1$ ), then an integration between 0 and 1 gives

$$\nu_k = \frac{\gamma}{k+1} + \int_0^1 x^k \psi(x+1) dx.$$

Thus, it follows from formula (2) that

$$\int_0^1 x^k \psi(x+1) dx = \sum_{j=0}^{k-1} (-1)^j \binom{k}{j} \zeta'(-j) + C_k \quad (k \geq 1).$$

**Remark 3** (Link with the Ramanujan summation: part I). Candelpergher et al. [11, Corollary 1] (see also [9, p. 82]) have established that

$$\sum_{n \geq 1}^{\mathcal{R}} H_n = \frac{3}{2}\gamma - \frac{1}{2} \ln 2\pi + \frac{1}{2},$$

and for any positive integer  $k$ ,

$$\sum_{n \geq 1}^{\mathcal{R}} n^k H_n = \left( \frac{1 - B_{k+1}}{k+1} \right) \gamma - \frac{1}{2} \ln 2\pi + \sum_{j=1}^k (-1)^j \binom{k}{j} \zeta'(-j) + R_k \quad \text{with } R_k \in \mathbb{Q},$$

where the symbol  $\sum^{\mathcal{R}}$  denotes the  $\mathcal{R}$ -sum of the series, i.e. the sum in the sense of Ramanujan's summation method [8, 9, 10, 11]. For an even integer  $k$ , we have  $B_{k+1} = 0$  and  $R_k = C_k + \frac{B_k}{2} - \frac{B_k}{2k}$ , then, in view of formula (2), these relations may be translated into the following identities:

$$\sum_{n \geq 1}^{\mathcal{R}} H_n = \nu_1 + \gamma - \frac{1}{2},$$

and for  $k \geq 1$ ,

$$\sum_{n \geq 1}^{\mathcal{R}} n^{2k} H_n = \nu_{2k} + \zeta'(-2k) + \frac{1 - 2k}{2} \zeta(1 - 2k) = \nu_{2k} + \zeta'(-2k) + (2k - 1) \frac{B_{2k}}{4k}. \quad (7)$$

In particular, we have

$$\sum_{n \geq 1}^{\mathcal{R}} n^2 H_n = \nu_2 + \zeta'(-2) + \frac{B_2}{4} = \nu_2 - \frac{\zeta(3)}{4\pi^2} + \frac{1}{24}.$$

## 2 The case $k = -1$

The case  $k = -1$  behaves differently from the previous case and must be studied separately. Denoting by  $\tau_1$  the constant  $\nu_{-1}$ , we have the following identities [12, p. 142]

$$\tau_1 = \int_0^1 \frac{\psi(x+1) + \gamma}{x} dx = \sum_{k=1}^{\infty} \frac{\ln(k+1)}{k(k+1)} = - \sum_{n=2}^{\infty} \zeta'(n) = 1.2577468869 \dots$$

Another interesting representation (communicated by I. V. Blagouchine) is

$$\tau_1 = \frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(3/2 + ix)}{(1/2 + ix) \cosh(\pi x)} dx.$$

Furthermore, we can write yet another relation that will be useful in the next section: for any integer  $p \geq 0$ , let  $\kappa_p$  be the constant defined by

$$\kappa_p := \sum_{n=1}^{\infty} \frac{|b_n|}{n^{p+1}},$$

where  $b_n$  are the *Bernoulli numbers of the second kind* defined by way of their generating function

$$\frac{x}{\log(x+1)} = \sum_{n=0}^{\infty} b_n x^n = 1 + \frac{x}{2} - \frac{x^2}{12} + \frac{x^3}{24} - \frac{19x^4}{720} + \frac{3x^5}{160} - \frac{863x^6}{60480} + \dots$$

The relation

$$g_n^{(1)} = \frac{1}{n!} \sum_{j=1}^n \frac{s(n,j)}{j+1} = b_n$$

is well-known [16, p. 267]. We have in particular,

$$\kappa_0 = \sum_{n=1}^{\infty} \frac{|b_n|}{n} = \gamma \quad \text{and} \quad \kappa_1 = \sum_{n=1}^{\infty} \frac{|b_n|}{n^2} = 0.5290529699\dots$$

The constants  $\kappa_1$  and  $\tau_1$  are linked by the relation [9, Eq. (3.23) p. 105]

$$\kappa_1 = \tau_1 + \gamma_1 + \frac{1}{2}\gamma^2 - \frac{1}{2}\zeta(2), \quad (8)$$

where  $\gamma_1$  denotes the first Stieljes constant [5, 8, 9].

$$\gamma_1 = \lim_{n \rightarrow \infty} \left\{ \sum_{j=1}^n \frac{\ln j}{j} - \frac{1}{2} \ln^2 n \right\} = -0.07281584548\dots$$

In terms of Ramanujan summation,  $\gamma_1$  is  $\sum_{n \geq 1}^{\mathcal{R}} \frac{\ln n}{n}$  (cf. [9, p. 67]), whereas  $\kappa_1$  is  $\sum_{n \geq 1}^{\mathcal{R}} \frac{H_n}{n}$  (cf. [9, Eq. (4.29) p. 133]).

### 3 Alternating series involving multiple zeta values

In this section, we consider a more general class of series of the previous type replacing zeta values with certain multiple zeta values. We prove our formula (3) and deduce some interesting consequences.

**Proposition 2.** For all integers  $p \geq 0$  and  $k \geq -1$ , let

$$\nu_{k,p} := \sum_{n=2}^{\infty} \frac{(-1)^n}{n+k} \zeta(n, \underbrace{1, \dots, 1}_p);$$

then

$$\nu_{k,p} = \sum_{n=1}^{\infty} \frac{|g_n^{(k+1)}|}{n^{p+1}},$$

where the rational numbers  $g_n^{(k)}$  are defined by equation (4). In particular, for  $p = 0$ , we have  $\nu_{k,0} = \nu_k$  and thus the identity

$$\nu_{k-1} = \sum_{n=1}^{\infty} \frac{|g_n^{(k)}|}{n} \quad (k \geq 0),$$

whereas for  $k = 0$ , we have

$$\nu_{0,p} = \sum_{n=1}^{\infty} \frac{|g_n^{(1)}|}{n^{p+1}} = \sum_{n=1}^{\infty} \frac{|b_n|}{n^{p+1}} = \kappa_p \quad (p \geq 0).$$

In order to prove Proposition 2, we begin by stating the following lemmas:

**Lemma 1.** For all integers  $j \geq 1$  and  $p \geq 0$ , we have

$$\int_0^1 \frac{\ln^j(1-x) \ln^p(x)}{x} dx = (-1)^{j+p} j! p! \zeta(j+1, \underbrace{1, \dots, 1}_p). \quad (9)$$

*Proof.* This follows directly from [19, Eqs. (2.27), (2.28)].  $\square$

**Lemma 2.** The Stirling numbers of the first kind  $s(n, j)$  with fixed  $j \geq 1$  admit the (vertical) exponential generating function [3, Eq. (2.8)]

$$\frac{\ln^j(1+x)}{j!} = \sum_{n=j}^{\infty} s(n, j) \frac{x^n}{n!} \quad (|x| < 1). \quad (10)$$

**Lemma 3.** For all integers  $n \geq 1$  and  $p \geq 0$ , we have

$$(-1)^p \int_0^1 x^{n-1} \ln^p(x) dx = \frac{p!}{n^{p+1}} \quad (11)$$

*Proof.* This is nothing else than [8, Eq. (41)] in the case where  $p$  is an integer.  $\square$

**Lemma 4.** For all integers  $n \geq 1$  and  $k \geq 0$ , we have

$$g_n^{(k)} = \frac{(-1)^{n+1}}{n!} \int_0^1 x^k (1-x)_{n-1} dx,$$

where  $(z)_n = z(z+1)(z+2) \cdots (z+n-1)$  is the Pochhammer symbol. In particular, this implies that

$$g_n^{(k)} = (-1)^{n+1} |g_n^{(k)}|. \quad (12)$$

*Proof.* Integration between 0 and 1 of the expansion

$$x^{k-1}x(x-1)\cdots(x-n+1) = \sum_{j=1}^n s(n, j)x^{j+k-1}$$

gives the required result.  $\square$

*Proof of Proposition 2.* Using successively formulas (9)–(12) above, we can write the following equalities:

$$\begin{aligned} \nu_{k,p} &= \sum_{j=1}^{\infty} \frac{(-1)^{j+1}}{j+k+1} \zeta(j+1, \underbrace{1, \dots, 1}_p) \\ &= \frac{(-1)^{p+1}}{p!} \sum_{j=1}^{\infty} \frac{1}{j+k+1} \int_0^1 \frac{\ln^j(1-x) \ln^p(x)}{j! x} dx \\ &= \frac{(-1)^{p+1}}{p!} \sum_{j=1}^{\infty} \frac{1}{j+k+1} \int_0^1 \left( \sum_{n=j}^{\infty} (-1)^n s(n, j) \frac{x^n}{n!} \right) \frac{\ln^p(x)}{x} dx \\ &= \frac{(-1)^{p+1}}{p!} \sum_{j=1}^{\infty} \frac{1}{j+k+1} \sum_{n=j}^{\infty} (-1)^n \frac{s(n, j)}{n!} \int_0^1 x^{n-1} \ln^p(x) dx \\ &= - \sum_{j=1}^{\infty} \frac{1}{j+k+1} \sum_{n=j}^{\infty} (-1)^n \frac{s(n, j)}{n! n^{p+1}} \\ &= - \sum_{n=1}^{\infty} \frac{(-1)^n}{n! n^{p+1}} \sum_{j=1}^n \frac{s(n, j)}{j+k+1} \\ &= \sum_{n=1}^{\infty} (-1)^{n+1} \left( \frac{1}{n!} \sum_{j=1}^n \frac{s(n, j)}{j+k+1} \right) \frac{1}{n^{p+1}} \\ &= \sum_{n=1}^{\infty} (-1)^{n+1} \frac{g_n^{(k+1)}}{n^{p+1}} = \sum_{n=1}^{\infty} \frac{|g_n^{(k+1)}|}{n^{p+1}}. \end{aligned}$$

This completes the proof.  $\square$

**Example 2.** For the first values of  $k \geq -1$ , we have the following expansions:

$$\tau_1 = \sum_{n=1}^{\infty} \frac{|g_n^{(0)}|}{n} = 1 + \frac{1}{8} + \frac{5}{108} + \frac{3}{128} + \frac{251}{18000} + \frac{95}{10368} + \dots,$$

$$\nu_0 = \sum_{n=1}^{\infty} \frac{|g_n^{(1)}|}{n} = \frac{1}{2} + \frac{1}{24} + \frac{1}{72} + \frac{19}{2880} + \frac{3}{800} + \frac{863}{362880} + \dots,$$

$$\nu_1 = \sum_{n=1}^{\infty} \frac{|g_n^{(2)}|}{n} = \frac{1}{3} + \frac{1}{48} + \frac{7}{1080} + \frac{17}{5760} + \frac{41}{25200} + \frac{731}{725760} + \dots,$$

$$\nu_2 = \sum_{n=1}^{\infty} \frac{|g_n^{(3)}|}{n} = \frac{1}{4} + \frac{1}{80} + \frac{1}{270} + \frac{11}{6720} + \frac{89}{100800} + \frac{5849}{10886400} + \dots.$$

**Corollary 1.** Let  $\zeta_H$  be the Apostol-Vu harmonic zeta function [4, 5, 6, 11] defined for  $\text{Re}(s) > 1$  by

$$\zeta_H(s) := \sum_{n=1}^{\infty} \frac{H_n}{n^s}.$$

We have the following elegant evaluation:

$$\sum_{n=2}^{\infty} (-1)^n \frac{\zeta_H(n)}{n} = \gamma_1 + \frac{1}{2}\gamma^2 + \frac{\pi^2}{12} = 0.916240149\dots \quad (13)$$

Another expression of this constant is  $\zeta''(0) + \frac{1}{2} \ln^2(2\pi) + \frac{\pi^2}{8}$  [5, Eq. (25.6.12)].

*Proof.* Using Proposition 2, we can write

$$\sum_{n=2}^{\infty} (-1)^n \frac{\zeta_H(n)}{n} = \nu_{0,1} - \tau_1 + \zeta(2) = \kappa_1 - \tau_1 + \zeta(2),$$

and thus, formula (13) is derived from equation (8).  $\square$

**Remark 4** (Link with the Ramanujan summation: part II). For  $s \in \mathbb{C}$ , let  $\zeta_H^{\mathcal{R}}$  be the function  $s \mapsto \sum_{n \geq 1}^{\mathcal{R}} H_n n^{-s}$  where  $\sum^{\mathcal{R}}$  stands for the Ramanujan summation. The function  $\zeta_H^{\mathcal{R}}$  is an entire function linked to the harmonic zeta function  $\zeta_H$  by the relation [11, Eq. (84)]

$$\zeta_H^{\mathcal{R}}(s) = \zeta_H(s) - \int_1^{\infty} x^{-s} (\psi(x+1) + \gamma) dx \quad \text{for } \text{Re}(s) > 1.$$

We have the identities

$$\zeta_H^{\mathcal{R}}(1) = \nu_{0,1} = \kappa_1, \quad \zeta_H^{\mathcal{R}}(0) = \nu_1 + \gamma - \frac{1}{2},$$

and formula (7) may be nicely rewritten

$$\zeta_H^{\mathcal{R}}(-2k) = \zeta_H(-2k) + \zeta'(-2k) + \nu_{2k}.$$

## 4 Link with the shifted Mascheroni series

Let us consider now the forward shifted Mascheroni series which are defined by

$$\sigma_r := \sum_{n=1}^{\infty} \frac{|b_{n+r}|}{n}, \quad \text{for } r = 0, 1, 2, \dots.$$

We have in particular  $\sigma_0 = \nu_0 = \gamma$ . The study of these series  $\sigma_r$  was the main subject of [13]. Among other things, we have established the following decomposition of  $\zeta'(-j)$  on the “basis” of  $\sigma_r$  [13, Proposition 3]:

$$\zeta'(-j) = \sum_{r=2}^{j+1} (-1)^{j-r} (r-1)! S(j, r-1) \sigma_r - \frac{B_{j+1}}{j+1} \gamma - \frac{B_{j+1}}{(j+1)^2}, \quad \text{for } j = 1, 2, 3, \dots,$$

where  $S(j, r)$  are Stirling numbers of the second kind; moreover, for  $j = 0$ , we also have a similar relation:

$$-\zeta'(0) = \frac{1}{2} \ln 2\pi = \sigma_1 + \frac{\gamma}{2} + \frac{1}{2}.$$

Then, substituting these relations into (2) enables us to write each series  $\nu_k$  with  $k \geq 1$  as an integral linear combination of  $\gamma, \sigma_1, \sigma_2, \dots, \sigma_k$  plus a rational number  $D_k$  which is closely linked to  $C_k$ . In this combination, the coefficient of  $\gamma$  is zero since it is equal to  $\frac{1}{k+1} \sum_{j=0}^k \binom{k+1}{j} B_j$  which vanishes by a well-known property of the Bernoulli numbers. Finally, equation (2) may be nicely rewritten in terms of  $\sigma_r$  as follows:

**Proposition 3.** For all integers  $k \geq 1$ , we have the relation

$$\nu_k = D_k - \sigma_1 + \sum_{r=2}^k (-1)^r (r-1)! \left( \sum_{j=r-1}^{k-1} \binom{k}{j} S(j, r-1) \right) \sigma_r \quad (14)$$

with

$$D_k = \frac{1}{k} + \sum_{j=1}^k \binom{k}{j} \frac{B_j H_j}{k+1-j}.$$

**Example 3.** For the first values of  $k$ , we have the following relations:

$$\begin{aligned} \nu_1 &= \frac{1}{2} - \sigma_1, \\ \nu_2 &= \frac{1}{4} - \sigma_1 + 2\sigma_2, \\ \nu_3 &= \frac{5}{24} - \sigma_1 + 6\sigma_2 - 6\sigma_3, \\ \nu_4 &= \frac{13}{72} - \sigma_1 + 14\sigma_2 - 36\sigma_3 + 24\sigma_4, \\ \nu_5 &= \frac{109}{720} - \sigma_1 + 30\sigma_2 - 150\sigma_3 + 240\sigma_4 - 120\sigma_5. \end{aligned}$$

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